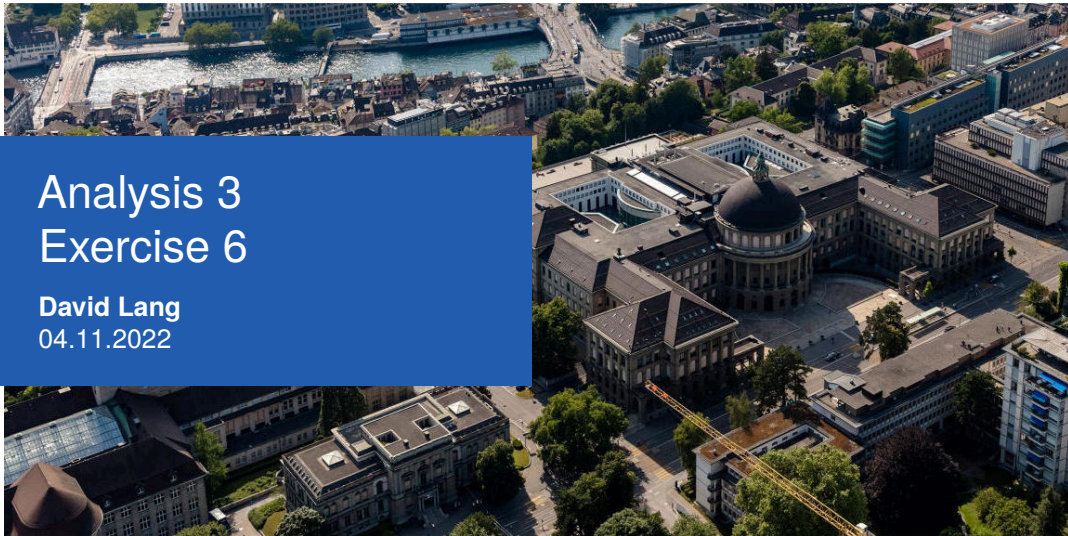


Analysis 3 Exercise 6

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04.11.2022



Outline

1. Serie 5 Review
2. Course Overview
3. Wave Equation
4. Canonical Form and Change of Variables
5. D'Alembert Formula
6. Examples
7. Tips for Serie 6

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Serie 5 Review

1. Weak solutions

- Not differentiable $\rightarrow$ no chain rule

2. Balance laws

- Characteristics are not straight lines anymore.

3. Multiple choice

- An antisymmetric matrix must have zeros on its diagonal. $A = -A^T$

4. Weak solutions II

- slope $1/c(u_0(s)) = 1/e^{-u_0(s)} = e^{u_0(s)}$

5. Finding shock waves

- $\frac{d}{dy}\gamma(y) = \gamma_y(y) = \gamma'(y)$

$$x = r(y)$$

$$\begin{cases} u(x,y) = \dots & x < r(y) \\ u(x,y) = \dots & x > r(y) \end{cases}$$

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Course Overview

- 1st order PDEs
 - Quasilinear first order PDEs
 - ▶ Method of characteristics
 - ▶ Conservation laws
- 2nd order PDEs
 - Hyperbolic PDEs
 - ▶ **Wave equation**
 - ▶ **D'Alembert formula**
 - ▶ Separation of variables
 - Parabolic PDEs
 - ▶ Heat equation
 - ▶ Maximum principle
 - ▶ Separation of variables
 - Elliptic PDEs
 - ▶ Laplace equation
 - ▶ Maximum principle
 - ▶ Separation of variables

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Wave Equation

The homogeneous wave equation in one (spatial) dimension has the form

$$u_{tt} - c^2 u_{xx} = 0, x \in \mathbb{R}, t > 0$$

$c \in \mathbb{R}$ is called the wave speed.

Note that $x \in \mathbb{R}$, which means that the problem can be thought of as the amplitude of the vibration of an **infinite** string.

This is the homogeneous wave equation, i.e. no external force.

If we impose boundary conditions (maybe only looking at $[0, L]$), then we will have to do some modifications, such as using the method of **Separation of Variables**.

$$au_{xx} + 2b u_{xy} + c u_{yy} + d u_x + e u_y + f u = g$$

$$a = -c^2 \quad c = 1$$

$$\Delta(u) = b^2 - ac = c^2 > 0 \quad \text{hyperbolic}$$

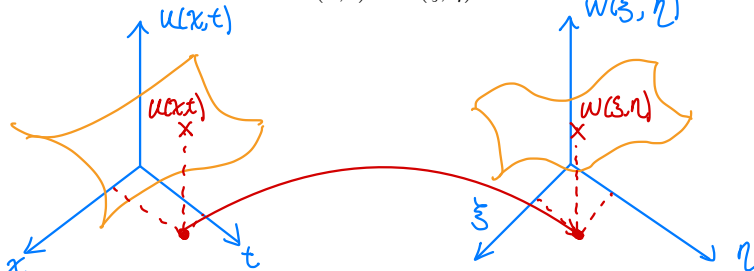
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7. Tips for Serie 6

Canonical Form and Change of Variables

$$\xi(x, t) = x + ct \quad \& \quad \eta(x, t) = x - ct$$

$$u(x, t) = w(\xi, \eta)$$



$$u_t = w_\xi \xi_t + w_\eta \eta_t \quad \& \quad u_x = w_\xi \xi_x + w_\eta \eta_x$$

$$u_{tt} = c^2(w_{\xi\xi} - 2w_{\xi\eta} + w_{\eta\eta}) \quad \& \quad u_{xx} = w_{\xi\xi} + 2w_{\xi\eta} + w_{\eta\eta}$$

$$u_{tt} - c^2 u_{xx} = 0 = -4c^2 w_{\xi\eta}$$

Canonical Form and Change of Variables

The result from the previous slide:

$$\frac{\partial}{\partial \eta} w_\xi = 0$$

w_ξ is independent of η

$$w_\xi(\xi, \eta) = f(\xi)$$

Integrate with respect to ξ , we get

$$w(\xi, \eta) = F(\xi) + G(\eta)$$

Transform back to the original coordinates

$$u(x, t) = F(x + ct) + G(x - ct)$$

backward *forward*

Characteristics of the Wave Equation

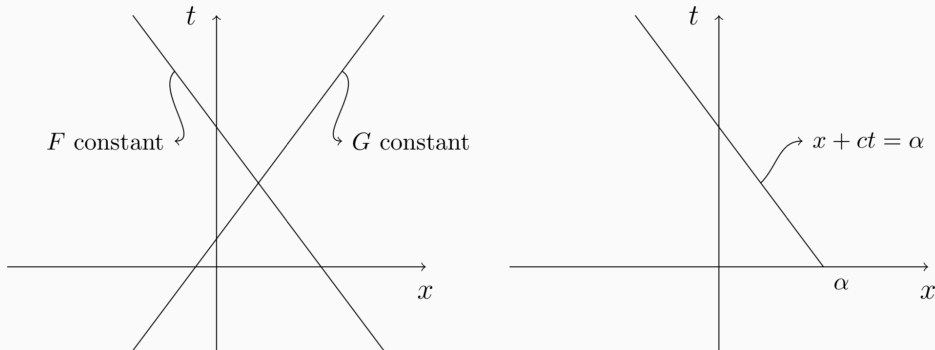
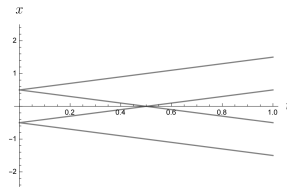
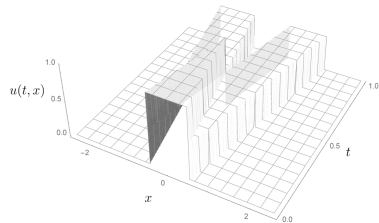
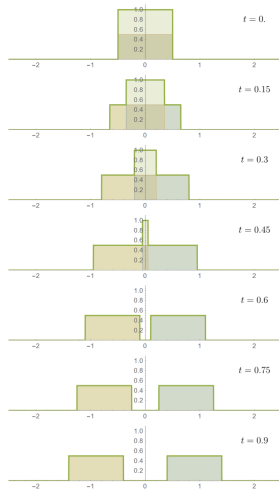


Figure 4.1: On the left, the characteristics where F and G are constant. On the right, the backward wave $F(x + ct)$.

Characteristics of the Wave Equation



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D'Alembert Formula

Cauchy problem for the wave equation:

$$\begin{cases} u_{tt} - c^2 u_{xx} = 0 & (x, t) \in \mathbb{R} \times (0, \infty) \\ u(x, 0) = f(x) \\ u_t(x, 0) = g(x) \end{cases}$$

We want $u(x, t)$ to have the form $F(x + ct) + G(x - ct)$.

By plugging in the values for $t = 0$, we could find out what F and G are and thus the general solution.

D'Alembert Formula

$$u(x, t) = \frac{f(x + ct) + f(x - ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(y) dy$$

Note: the value of the solution at (x_0, t_0) is only influenced by the values of f and g in $[x_0 - ct_0, x_0 + ct_0]$

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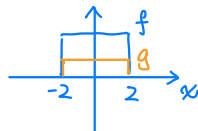
Example 1

Let $u(x, t)$ be the solution of the following initial value problem

$$\begin{cases} u_{tt} = 4u_{xx} & x \in \mathbb{R}, t > 0 \\ u(x, 0) = f(x) & x \in \mathbb{R} \\ u_t(x, 0) = g(x) & x \in \mathbb{R} \end{cases}$$

$$f(x) = \begin{cases} 3 & |x| \leq 2 \\ 0 & |x| > 2 \end{cases}$$

$$g(x) = \begin{cases} 1 & |x| \leq 2 \\ 0 & |x| > 2 \end{cases}$$



Find $u(1, 1)$

$$u(x, t) = \frac{f(x+ct) + f(x-ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(y) dy$$

$$c = 2$$

$$\begin{aligned} u(1, 1) &= \frac{f(1+2) + f(1-2)}{2} + \frac{1}{2 \cdot 2} \int_{1-2}^{1+2} g(y) dy \\ &= \frac{f(3) + f(-1)}{2} + \frac{1}{4} \int_{-1}^3 1 dy \end{aligned}$$

$$\begin{aligned} &= \frac{0+3}{2} + \frac{1}{4} \int_{-1}^3 1 dy \\ &= \frac{3}{2} + \frac{1}{4} \cdot 3 \\ &= \frac{9}{4} \end{aligned}$$

Example 1

Let $u(x, t)$ be the solution of the following initial value problem

$$\begin{cases} u_{tt} = 4u_{xx} & x \in \mathbb{R}, t > 0 \\ u(x, 0) = f(x) & x \in \mathbb{R} \\ u_t(x, 0) = g(x) & x \in \mathbb{R} \end{cases}$$

$$f(x) = \begin{cases} 3 & |x| \leq 2 \\ 0 & |x| > 2 \end{cases}$$

$$g(x) = \begin{cases} 1 & |x| \leq 2 \\ 0 & |x| > 2 \end{cases}$$

Find $\lim_{t \rightarrow \infty} u(1, t)$

$$\begin{aligned} \lim_{t \rightarrow \infty} u(1, t) &= \lim_{t \rightarrow \infty} \frac{1}{2} [f(1+2t) + f(1-2t)] + \frac{1}{4} \int_{1-2t}^{1+2t} g(y) dy \\ &= 0 + \frac{1}{4} \int_{-2}^2 1 dy \\ &= 0 + \frac{1}{4} \cdot 4 \\ &= 1 \end{aligned}$$

Example 2

Consider the initial value problem with zero boundary condition

$$\begin{cases} u_{tt} - u_{xx} = 0 & (x, t) \in (0, \infty) \times (0, \infty) \\ u(0, t) = 0 & t \in (0, \infty) \\ u(x, 0) = x^4 & x \in [0, \infty) \\ u_t(x, 0) = \sin(x) & x \in [0, \infty) \end{cases}$$

Evaluate $u(2, 1)$ and $u(1, 2)$. In which of the two points $((2, 1)$ or $(1, 2))$ is the solution unaffected by the boundary condition at $x = 0$?

If $v(x, t)$ is an odd function, then it automatically satisfies $v(0, t) = 0$ for all t .
Extend $u(x, 0)$ & $u_t(x, 0)$ as odd functions to all $\mathbb{R}$.

$$\begin{cases} v_{tt} - v_{xx} = 0 & (x, t) \in \mathbb{R} \times (0, \infty) \\ v(x, 0) = x^3 \cdot |x| & x \in \mathbb{R} \\ v_t(x, 0) = \sin(x) & x \in \mathbb{R} \end{cases}$$

Example 2

Consider the initial value problem with zero boundary condition

$$\begin{cases} u_{tt} - u_{xx} = 0 & (x, t) \in (0, \infty) \times (0, \infty) \\ u(0, t) = 0 & t \in (0, \infty) \\ u(x, 0) = x^4 & x \in [0, \infty) \\ u_t(x, 0) = \sin(x) & x \in [0, \infty) \end{cases}$$

The D'Alembert formula is derived for $x \in \mathbb{R}$, we need to modify the problem before applying it.

We should define a new problem on $\mathbb{R}$ so that its solution provides the correct result if we just focus on $x \geq 0$.

Example 2

$$\begin{cases} u_{tt} - u_{xx} = 0 & (x, t) \in (0, \infty) \times (0, \infty) \\ u(0, t) = 0 & t \in (0, \infty) \\ u(x, 0) = x^4 & x \in [0, \infty) \\ u_t(x, 0) = \sin(x) & x \in [0, \infty) \end{cases} \quad \begin{cases} v_{tt} - v_{xx} = 0 & (x, t) \in \mathbb{R} \times (0, \infty) \\ v(x, 0) = x^3 \cdot |x| & x \in \mathbb{R} \\ v_t(x, 0) = \sin(x) & x \in \mathbb{R} \end{cases}$$

We get $u(x, t)$ by just restricting $v(x, t)$ to $x \geq 0$.

$$v(x, t) = \frac{1}{2} [(x+t)^3 |x+t| + (x-t)^3 |x-t|] + \frac{1}{2} \int_{x-t}^{x+t} \sin(y) dy$$

$$= \frac{1}{2} [(x+t)^3 |x+t| + (x-t)^3 |x-t|] - \frac{1}{2} [\cos(x+t) - \cos(x-t)]$$

$$v(1, 2) = \frac{1}{2} (3^4 - 1^4) + \frac{1}{2} (\cos(1) - \cos(3)) = 40 + \frac{1}{2} (\cos(1) - \cos(3)) \quad u(1, 2) = v(1, 2) \quad x=1 > 0$$

$$v(2, 1) = \frac{1}{2} (3^4 + 1^4) + \frac{1}{2} (\cos(1) - \cos(3)) = 41 + \frac{1}{2} (\cos(1) - \cos(3)) \quad u(2, 1) = v(2, 1) \quad x=2 > 0$$

domain of dependence:

$(1, 2): [x_0 - ct_0, x_0 + ct_0] = [-1, 3] \rightarrow$ affected by the zero boundary condition

$(2, 1): [x_0 - ct_0, x_0 + ct_0] = [1, 3] \rightarrow$ unaffected by the zero boundary condition

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Tips for Serie 6

1. Wave equation

- D'Alembert's Formula

$$\sin(a+b) + \sin(a-b) = 2\sin(a)\cos(b) \text{ \& } \cos(a+b) + \cos(a-b) = 2\cos(a)\cos(b)$$

2. Wave equation's anatomy

- (a) Chapter 4.2 in Script: The Cauchy problem and d'Alembert's formula

- (b) Apply d'Alembert's formula directly

$$\sin(a+b) + \sin(a-b) = 2\sin(a)\cos(b) \text{ \& } \cos(a+b) + \cos(a-b) = 2\cos(a)\cos(b)$$

3. Propagation of symmetries from initial data

- (a) Chapter 4.5 in Script: Symmetry of the wave equation

- (b) Periodic $\rightarrow$ Fourier series, even $\rightarrow$ Which terms of the Fourier Series disappear?

4. Multiple choice

-

5. Time reversible

- Check the properties one by one.

6. Zero boundary condition

- Example 2 of today's exercise.

Before the next lecture:

1. 3Blue1Brown: But what is a partial differential equation?
<https://youtu.be/ly4S0oi3Yz8>
2. 3Blue1Brown: Solving the Heat Equation
<https://youtu.be/TolXSwZ1pJU>

References:

1. Lecture notes on the course website.
2. “An Introduction to Partial Differential Equations” by Yehuda Pinchover and Jacob Rubinstein
3. NDSU lecture notes